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New Insights on Particle Acceleration at Non-relativistic Shocks

Damiano Caprioli University of Chicago



Collisionless Shocks

Mediated by collective electromagnetic interactions Show prominent non-thermal activity

Now studied in laboratory with laser experiments!





A universal acceleration mechanism

Fermi mechanism (Fermi, 1949): random elastic collisions lead to energy gain

PHYSICAL REVIEW

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On the Origin of the Cosmic Radiation

ENRICO FERMI Institute for Nuclear Studies, University of Chicago, Chicago, Illinois (Received January 3, 1949)

A theory of the origin of cosmic radiation is proposed according to which cosmic rays are originated and accelerated primarily in the interstellar space of the galaxy by collisions against moving magmetic fields. One of the features of the theory is that it yields naturally an inverse power law for the spectral distribution of the cosmic rays. The chief difficulty is that it fails to explain in a straightforward way the heavy nuclei observed in the primary radiation.



APRIL 15, 1949



Astroplasmas from first principles

Full-PIC approach Ø Define electromagnetic fields on a grid Move particles via Lorentz force Second Evolve fields via Maxwell equations B Computationally very challenging!

A Hybrid approach: Fluid electrons - Kinetic protons (Winske & Omidi; Burgess et al., Lipatov 2002; Giacalone et al. 1993,1997,2004-2013; DC & Spitkovsky 2013-2015,...)

massless electrons for more macroscopical time/length scales 0

$$\mathbf{E} = -\frac{\mathbf{V}_i}{c} \times \mathbf{B} + \frac{1}{4\pi n e} \left(\nabla \times \mathbf{B}\right) \times \mathbf{B} - \frac{T_i}{n}$$



B













dHybrid code (Gargaté et al, 2007; DC & Spitkovsky 2014, Haggerty & DC, 2018)



 $x [c / \omega_{pi}]$





Spectrum evolution

Acceleration efficiency: ~15% of the shock bulk energy!



DC & Spitkovsky, 2014a

• Diffusive Shock Acceleration: non-thermal tail with universal spectrum $f(p) \propto p^{-4}$





CR-driven Magnetic-Field Amplification



Initial B field M_s=M_A=30

DC & Spitkovsky, 2013

$$\begin{array}{c|c} 4000 & 5000 & 6000 & 7000 & 8000 \\ \hline x[c/\omega_p] \\ B_z & (t=2\omega_c^{-1}) \\ \hline 4000 & 5000 & 6000 & 7000 & 8000 \\ \hline x[c/\omega_p] \end{array}$$



3D simulations of a parallel shock



DC & Spitkovsky, 2014a







Parallel vs Oblique shocks













Dependence on shock strength (M_A) and inclination





More B amplification for stronger (higher M_A) shocks

Output Different flavors of CR-driven streaming instabilities (Amato & Blasi 2009; DC & Spitkovsky 2014b)

Study how CRs diffuse in the self-generated turbulence Bohm-like diffusion (DC & Spitkovsky 2014c)







SN 1006: a parallel accelerator



X-ray emission: red=thermal white=synchrotron





B amplification and ion acceleration where the shock is parallel



DC & Spitkovsky, 2014a

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Cosmic Wealth

IF U.S. LAND MASS WERE DIVIDED LIKE U.S. WEALTH

OWN THIS 9% WOULD OWN THIS

WOULD OWN

THIS RED DOT

30% WOULD OWN THIS

20% WOULD OWN THIS

1% WOULD

Source: Wikipedia

The top 1% carries ~one third of the total US wealth

deepened.









How to Become Non-Thermal: the Injection Problem

What determines the fraction of particles that become CRs?





 $x [c / \omega_{pi}]$

3000









Particle Injection - Simulations



Fraction of

Time $t = 80.100 \omega_c^{-1}$



DC, Pop & Spitkovsky, 2015





Encounter with the shock barrier

Low barrier (reformation)

average $|e\Delta\Phi|$



lons advected downstream, and thermalized

To overrun the shock, ions need a minimum Eini, increasing with 8 (DC, Pop & Spitkovsky 15) Ion fate determined by barrier duty cycle (~25%) and shock inclination After N SDA cycles, only a fraction $\eta \sim 0.25^{N}$ has not been advected \otimes For $\vartheta = 45^{\circ}$, $E_{ini} \sim 10E_0$, which requires N~3 -> $\eta \sim 1\%$ \odot For ϑ >45°, E_{ini} >10 E_0 , hence N>3 and η <<1%

High barrier (overshoot)

$|e\Delta\Phi| > mV_x^2/2$



lons reflected upstream, and energized via Shock Drift Acceleration







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Proton DSA: Summary

Shock Acceleration can be efficient
CRs amplify B via streaming instability
DSA efficient at parallel, strong shocks
Injection via specular reflection and shock-drift acceleration

What about other ions?







What if there are already energetic particles (seeds)?

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Diffusive Shock Re-Acceleration

 \otimes $\vartheta = 60^{\circ}$ shock with *isotropic* seeds with $E_{CR} = 3E_{sh}$; $n_{CR} = 0.01$ (DC, Zhang, Spitkovsky, 2018)

Seeds are effectively reflected at the shock, amplify the upstream B, and undergo DSA: DSRA!





Efficiency

80



 \odot Seed DSRA independent of ϑ , about 4x the initial CR energy density Absolute efficiency depends on seed energy density Also electrons can be reaccelerated!

 \oslash (45°< ϑ <70°): Boosted to few % $OC(\sqrt[9]{70^{\circ}})$: No proton DSA



Quasi-Perpendicular SEEDED Shocks

 \otimes $\vartheta = 80^{\circ}$ quasi-perp shock with seeds $E_{CR} = 3E_{sh}$ Seeds diffuse but their spectrum is steeper than DSA Non-thermal protons only downstream









The Current in Reflected CRs

\circ Depends on the fraction of reflected seeds, n, and their speed, v_r







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A Universal Current in Reflected CRs



 $\circ \eta$ and v_r "magically" balance their dependence on ϑ and M exactly: $J_{CR} = en_{CR}V_{sh}$ Easy explanation: CRs tend to become isotropic at the shock, in the shock frame: they become anisotropic in the upstream frame For SNRs and Galactic CRs: T_{stream inst}~10yr

Minimum level of B-amplification for shocks in the ISM



What about electrons?

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Ion vs Electron Injection

Ions injected by specular reflection Their magnetic moment $W_{\perp} = p_{\perp}^2/B$ is not conserved: the shock is evolving on their gyro-time! Electrons cannot be reflected by the shock potential barrier, but conserve their W \odot ∇ B-drift + shock drift acceleration $\frac{W_{\perp}}{qB} \frac{\mathbf{B} \times \nabla B}{B^2}$ $\mathbf{E} = -\mathbf{V}_1/c imes \mathbf{B}$ $\mathbf{v}_{\nabla B} =$ Electron injection requires oblique shocks! How can we have simultaneous acceleration of ions and electrons?









Electron Acceleration

\diamond Full PIC simulations (Tristan-MP code) M=20, V_{sh}=0.1c, quasi-parallel (ϑ =30°) 1D shock





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Section pre-heating via SDA





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What is the feedback of CRs on SNR evolution?

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SNR Evolution in a Thin-Shell

Ejecta-dominated stage: R_{sh}~V
Sedov-Taylor (adiabatic) stage:
Radiative stage (T_{sh}<~10⁶K)
Pressure-driven snowplow (P_{ho}
Momentum-driven snowplow

SNRs deposit energy and momentum in the interstellar medium Crucial for feedback that can suppress star formation!

/sht		
R _{sh} ~t 2/5		VIsh
$rac{P_0}{D_t}$	Phot	P ₀
(Phot~Po)	$\frac{d(M_{\rm sh}v_{\rm SNR})}{dt} = 4\pi r_{\rm SNR}^2 (P_{\rm ho})$	ot — 1







