

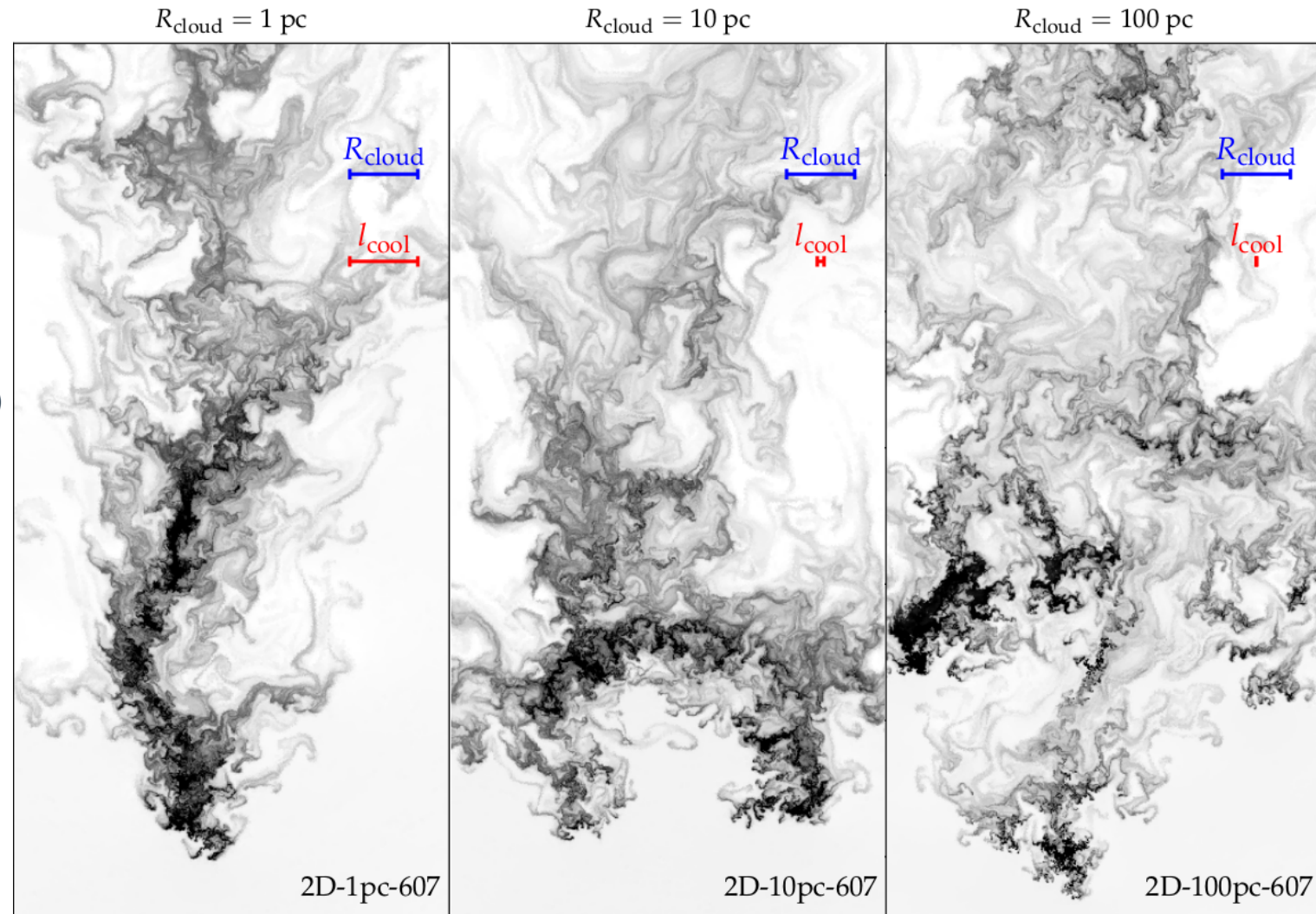
Shocks in 2D and 3D: Implications for shattering of accelerated gas clouds

Martin Sparre

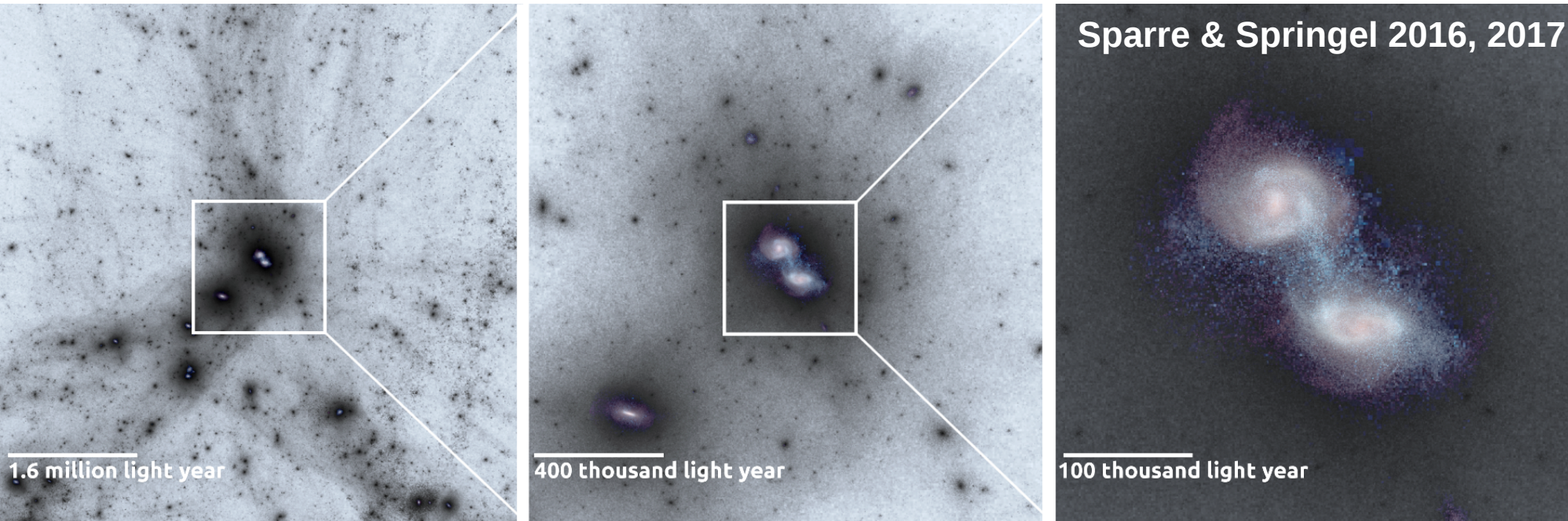
Potsdam Uni. / AIP

Collaborators:

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Mark Vogelsberger (MIT),
Volker Springel (HITS/MPA)



Cosmological simulations



Cosmological simulations resolve the cosmological large-scale structure and the 50-100 pc structure of galaxies.



A characteristic scale for cold gas

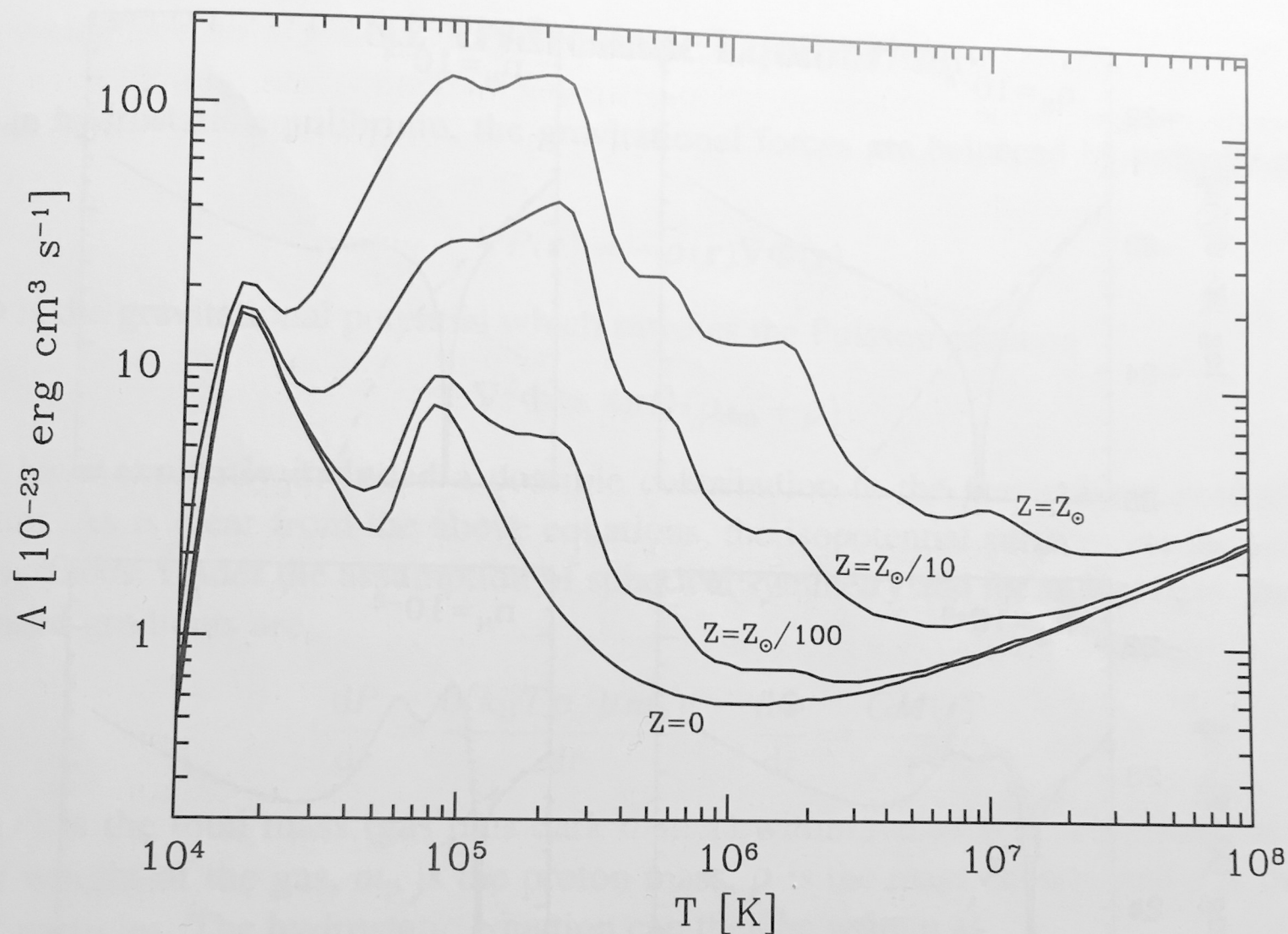
Michael McCourt,^{1★†} S. Peng Oh,^{1★} Ryan O’Leary² and Ann-Marie Madigan^{2,3}

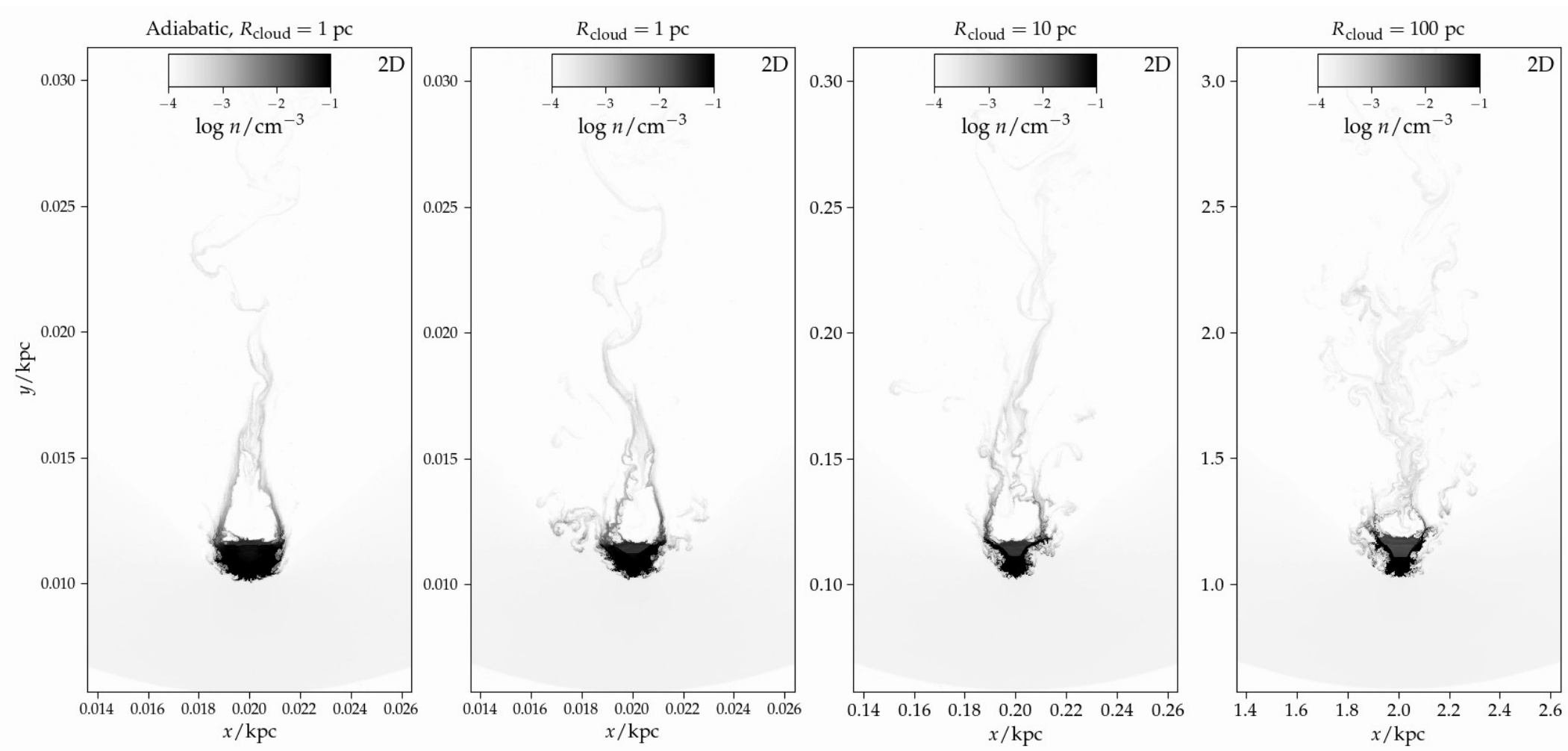
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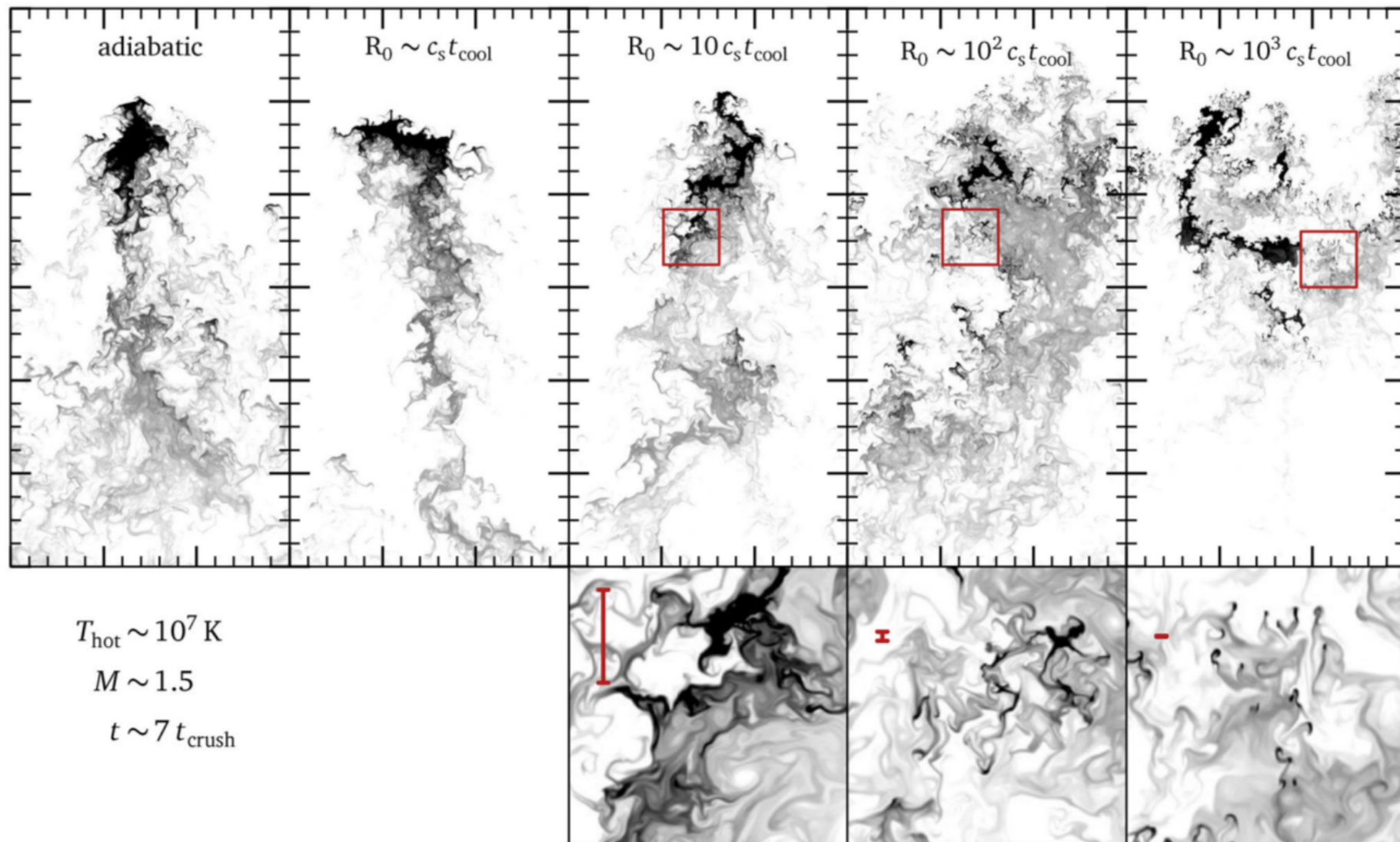
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A cloud interacting with a hot wind

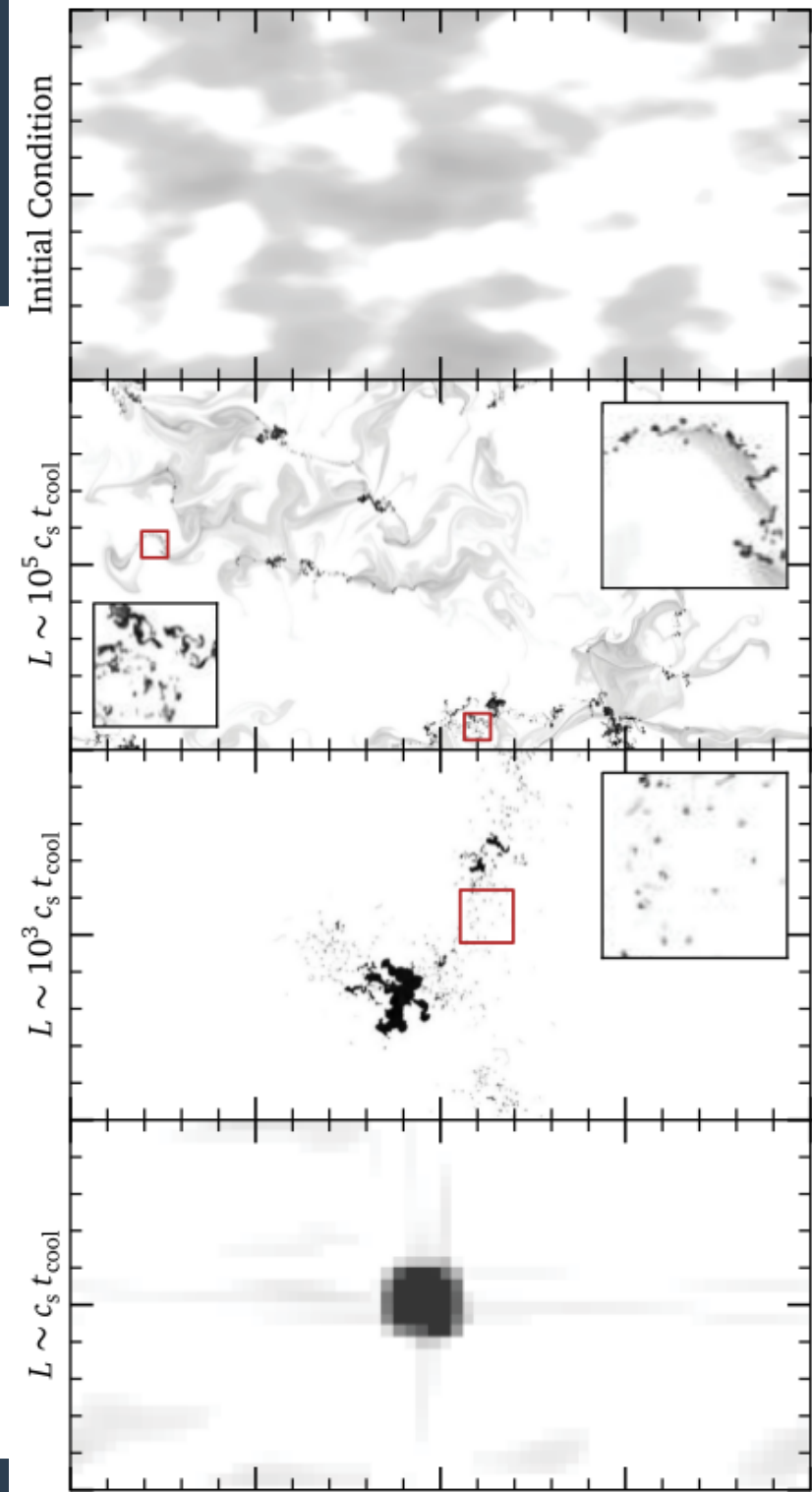


Take-away messages from McCourt:

- Clouds larger than $c_s t_{\text{cool}}$ “shatter” into smaller clouds.
- Gas in (the outer parts) of galaxy might be in a different physical state than revealed by cosmological simulations.

These results are based on 2D simulations. My work:

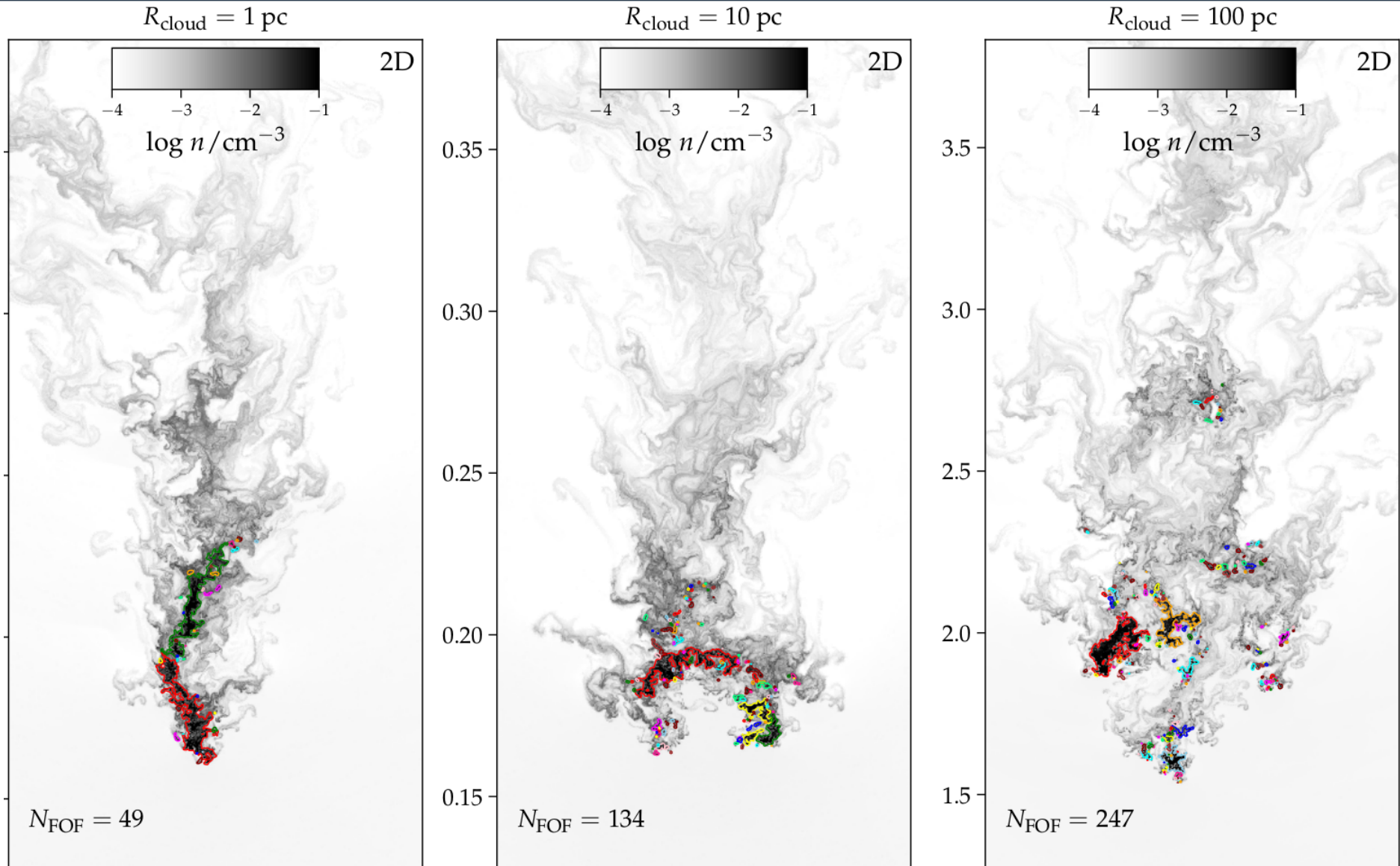
- 1) quantify fragmentation
- 2) 3D runs



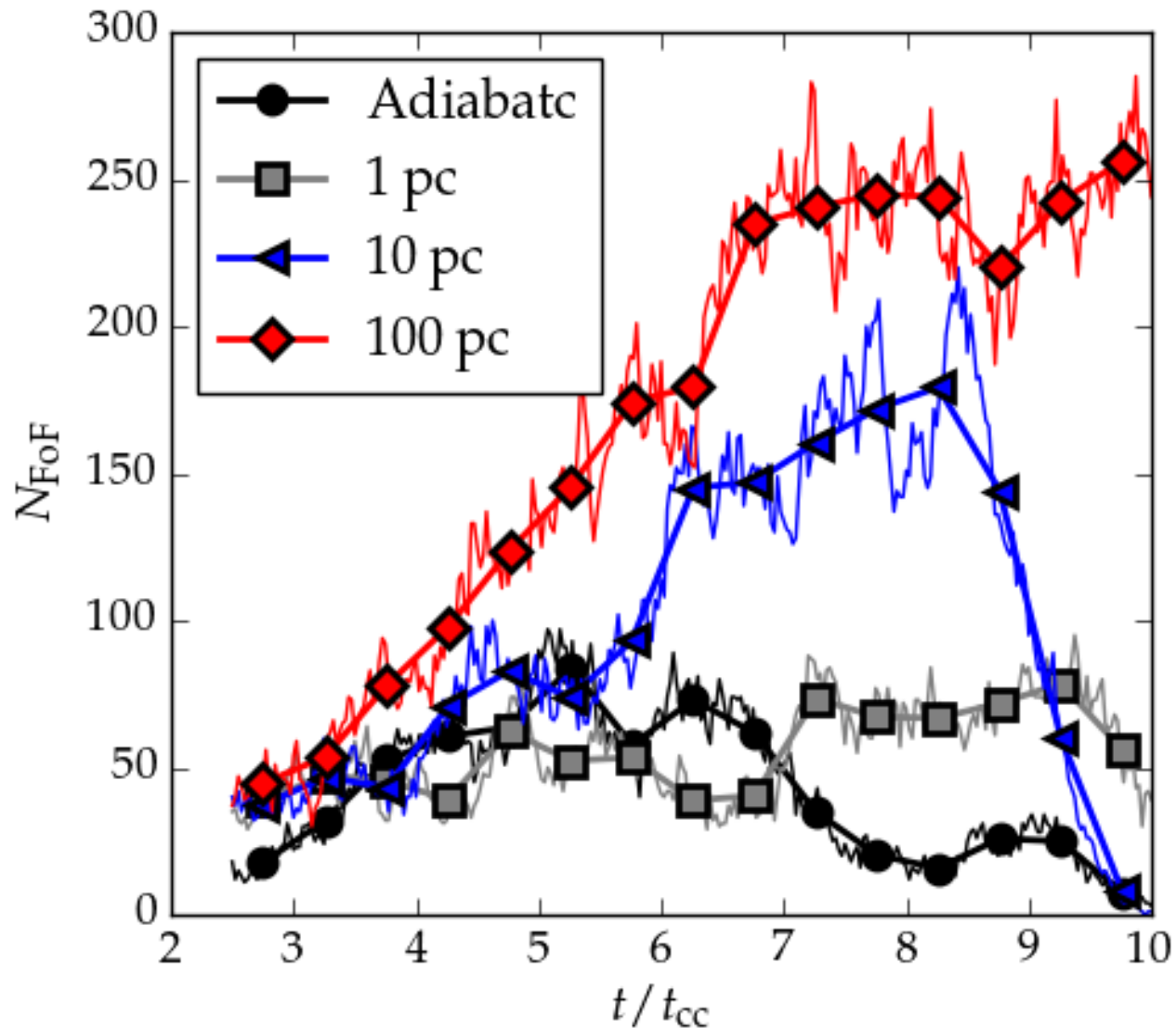
Quantifying the fragmentation



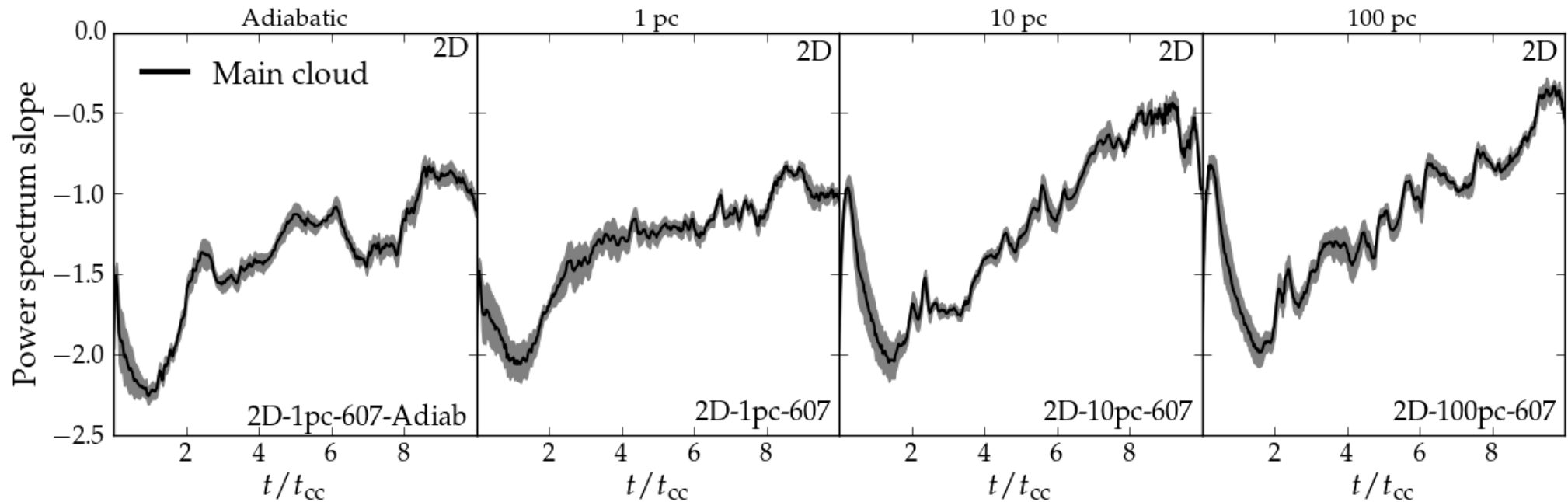
Friends-of-Friends analysis



Clouds larger than $c_s t_{\text{cool}}$ fragment



A shallow power spectrum is seen for clouds larger than $c_s t_{\text{cool}}$



- We have reproduced the McCourt results, and quantified the fragmentation.

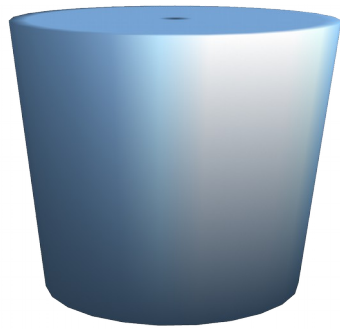


Two vs. three dimensions

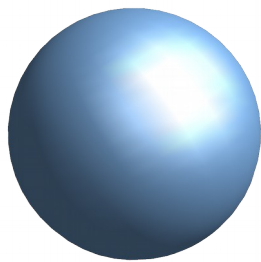


Potential flow solutions

A *potential flow* with $\text{div } \mathbf{v} = 0$ and $\text{curl } \mathbf{v} = \mathbf{0}$



$$\nabla \cdot \mathbf{v} = \frac{1}{\rho} \frac{\partial [\rho v_\rho]}{\partial \rho} + \frac{1}{\rho} \frac{\partial v_\phi}{\partial \phi},$$
$$\nabla \times \mathbf{v} = \frac{1}{\rho} \left[\frac{\partial [\rho v_\phi]}{\partial \rho} - \frac{\partial v_\rho}{\partial \phi} \right] \hat{\mathbf{z}},$$

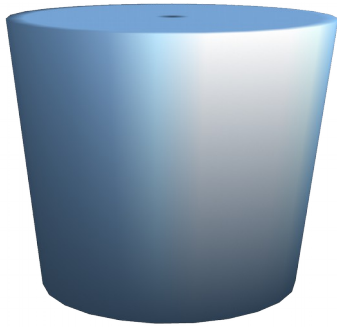


$$\nabla \cdot \mathbf{v} = \frac{1}{r^2} \frac{\partial [r^2 v_r]}{\partial r} + \frac{1}{r \sin \theta} \frac{\partial [v_\theta \sin \theta]}{\partial \theta},$$
$$\nabla \times \mathbf{v} = \frac{1}{r} \left[\frac{\partial [r v_\theta]}{\partial r} - \frac{\partial v_r}{\partial \theta} \right] \hat{\phi},$$



Potential flow solutions

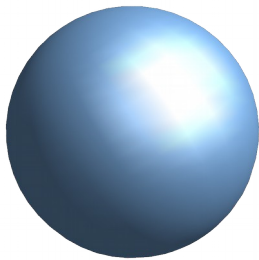
A *potential flow* has $\text{div } \mathbf{v} = 0$ and $\text{curl } \mathbf{v} = 0$



$$v_\rho = v_{\text{inject}} \left(1 - \frac{R^2}{\rho^2} \right) \cos \phi,$$

$$v_\phi = -v_{\text{inject}} \left(1 + \frac{R^2}{\rho^2} \right) \sin \phi,$$

$$v_z = 0$$



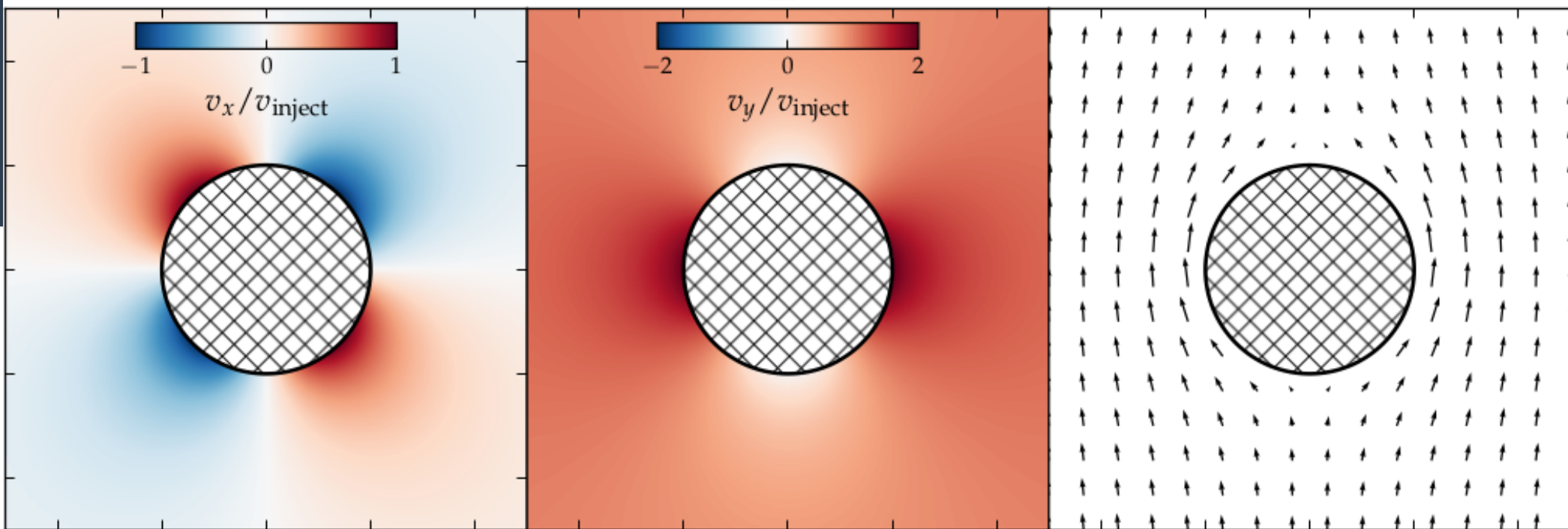
$$v_r = v_{\text{inject}} \left(1 - \frac{R^3}{r^3} \right) \cos \theta,$$

$$v_\theta = -v_{\text{inject}} \left(1 + \frac{R^3}{2r^3} \right) \sin \theta,$$

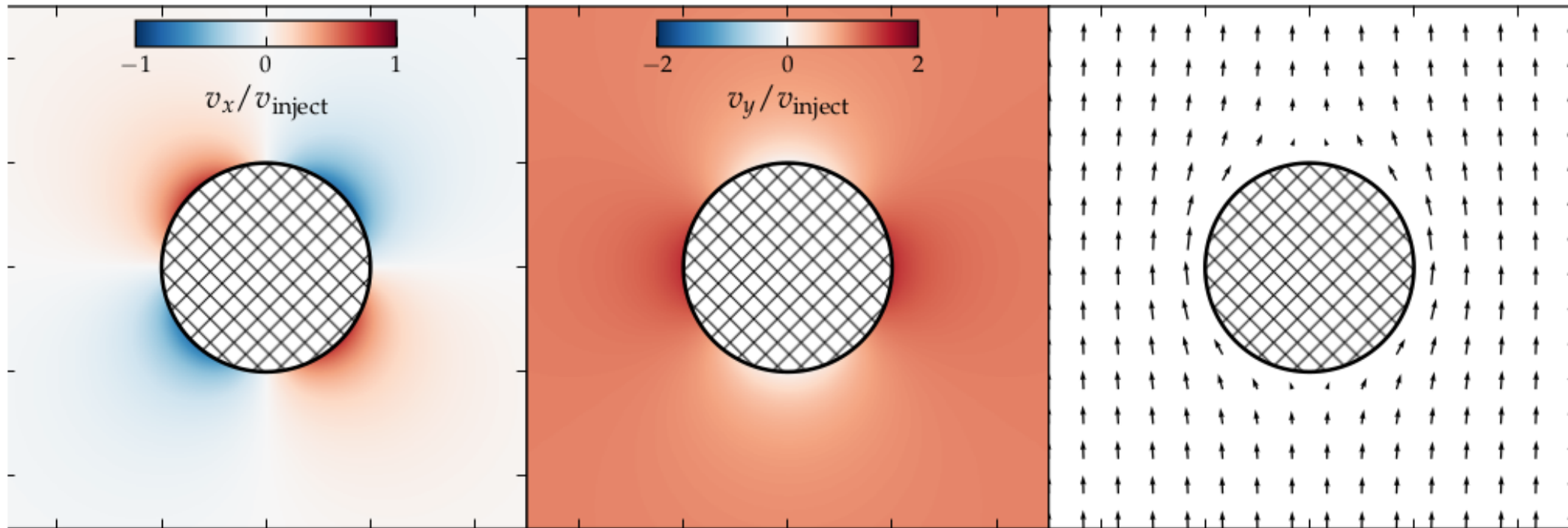
$$v_\varphi = 0,$$



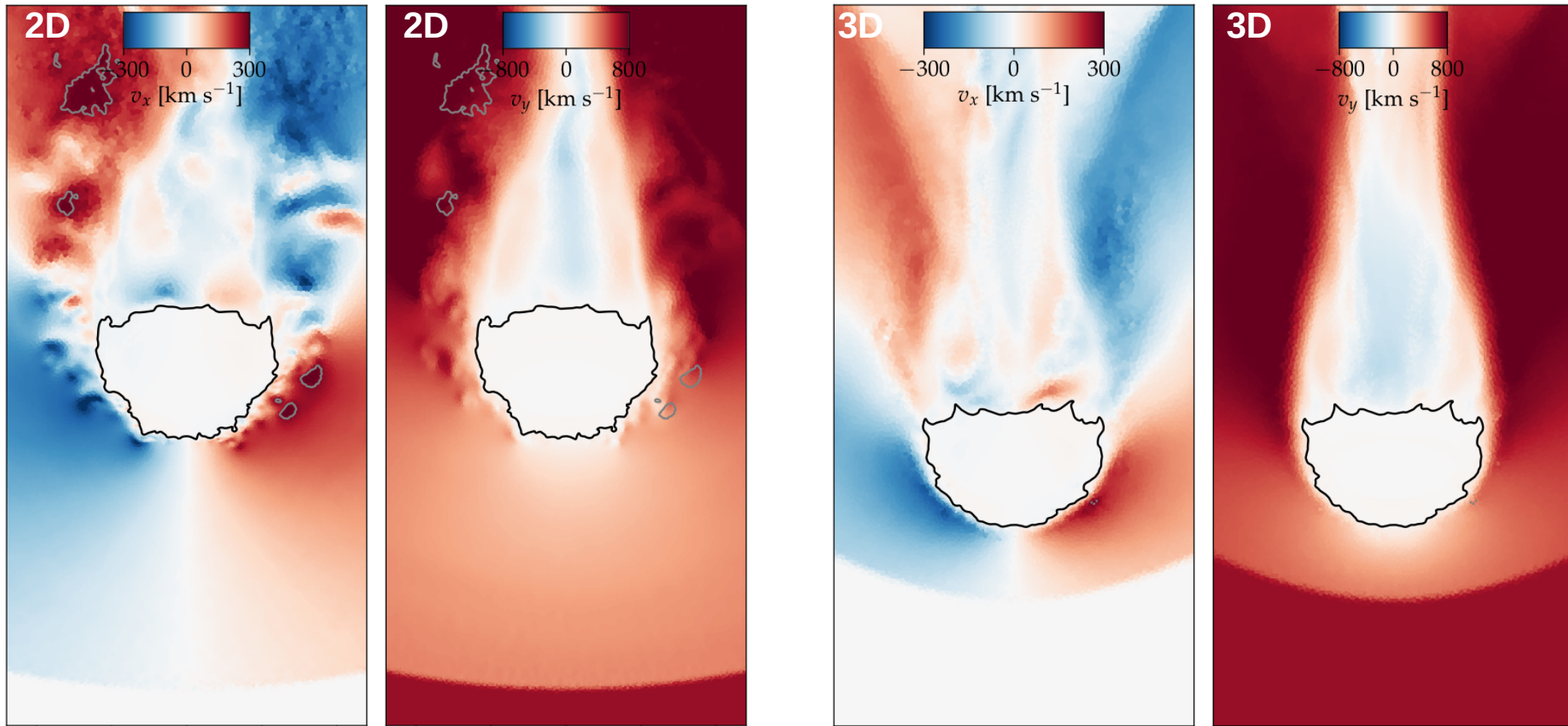
2D



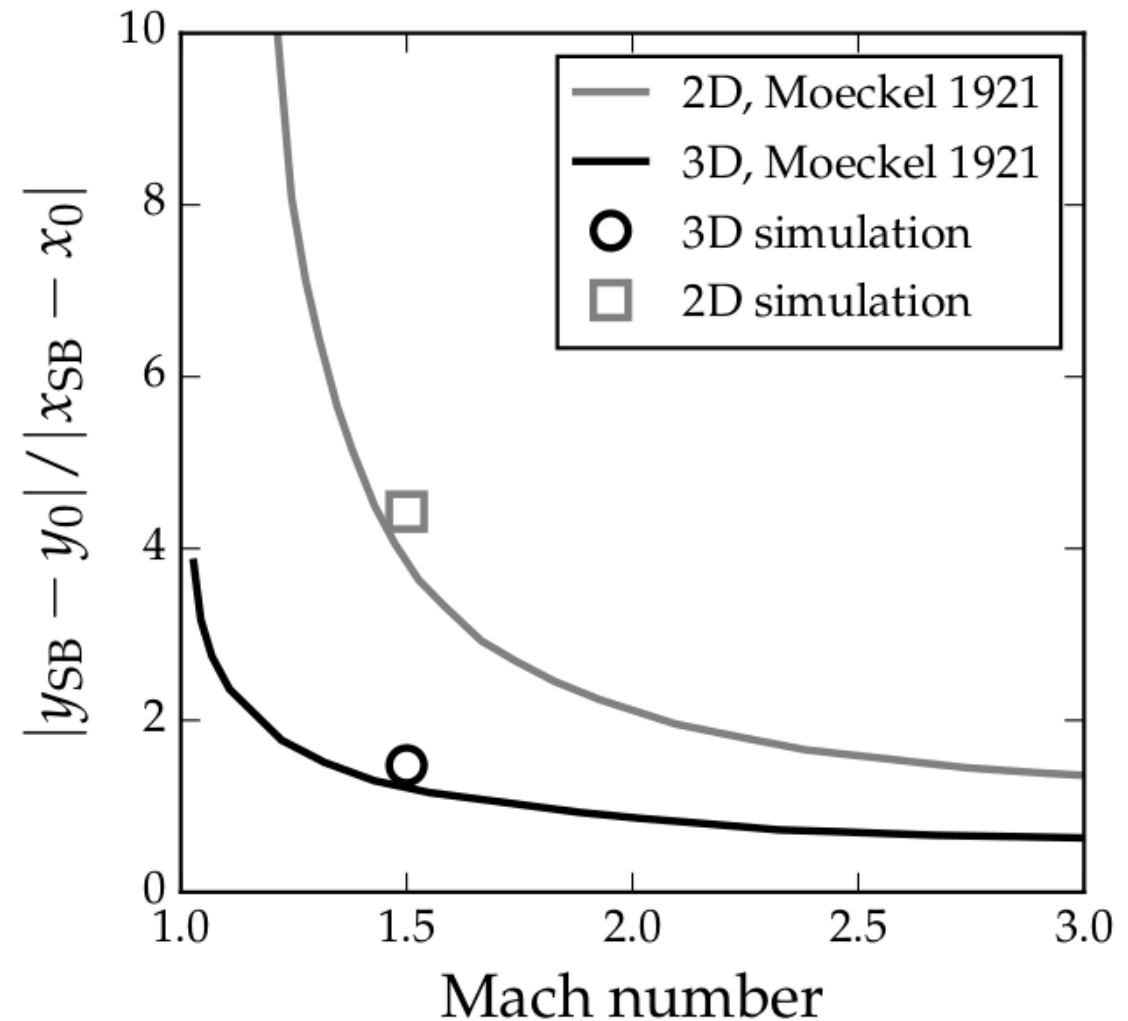
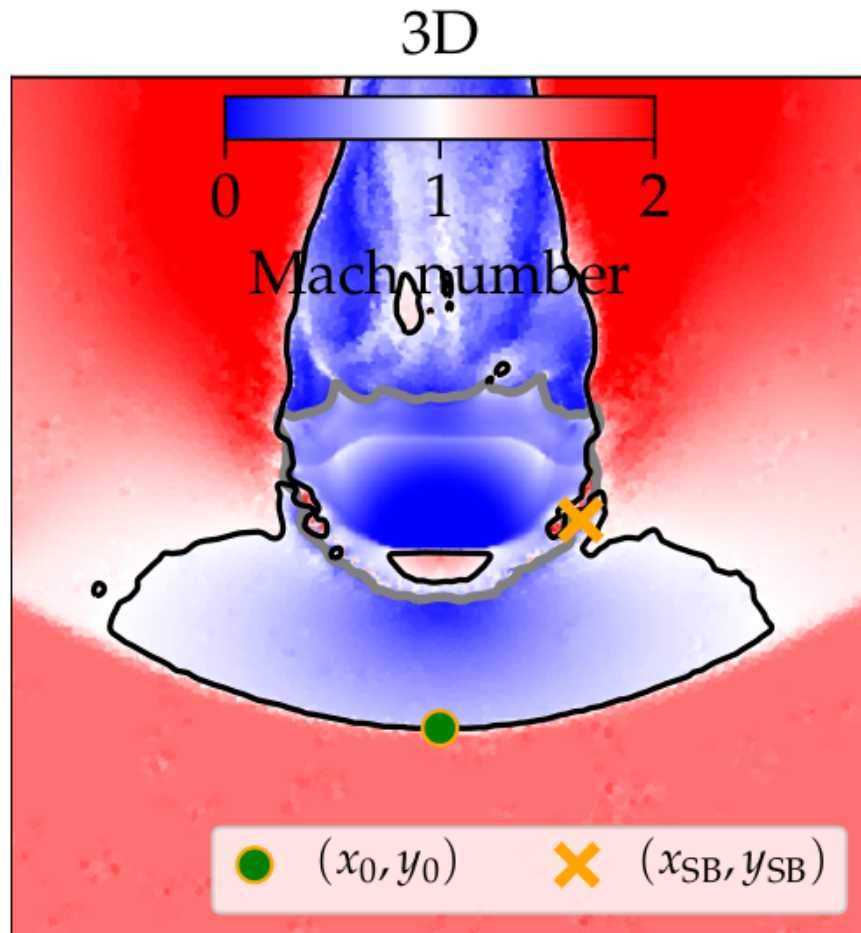
3D



The velocity field in 2D and 3D



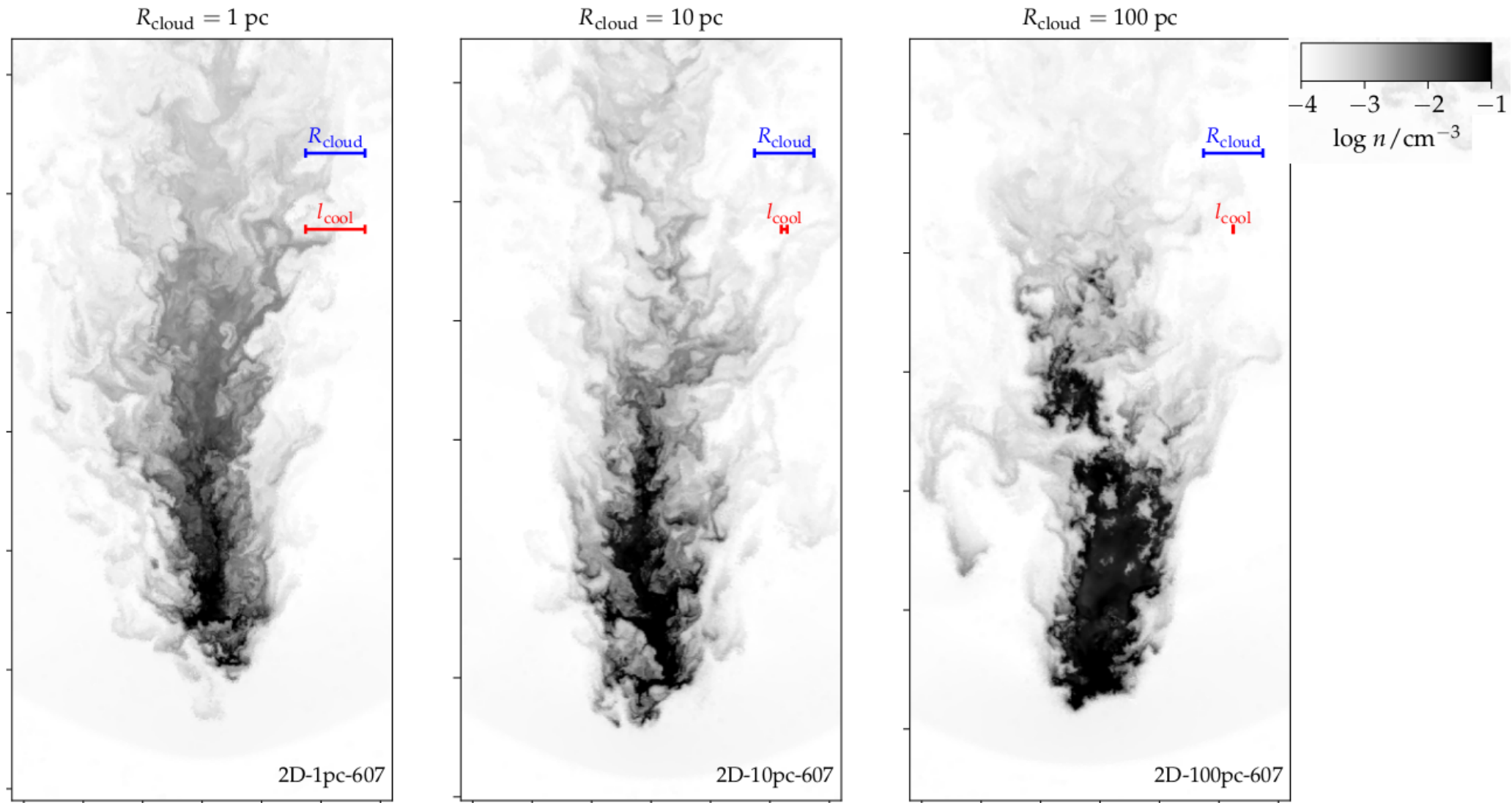
The standoff distance in 2D and 3D



- **Standoff-distance is larger in 2D than in 3D**
 - **Velocity around the cloud is largest in 2D.**
- 3D is necessary**



Shattering also occur in 3D!

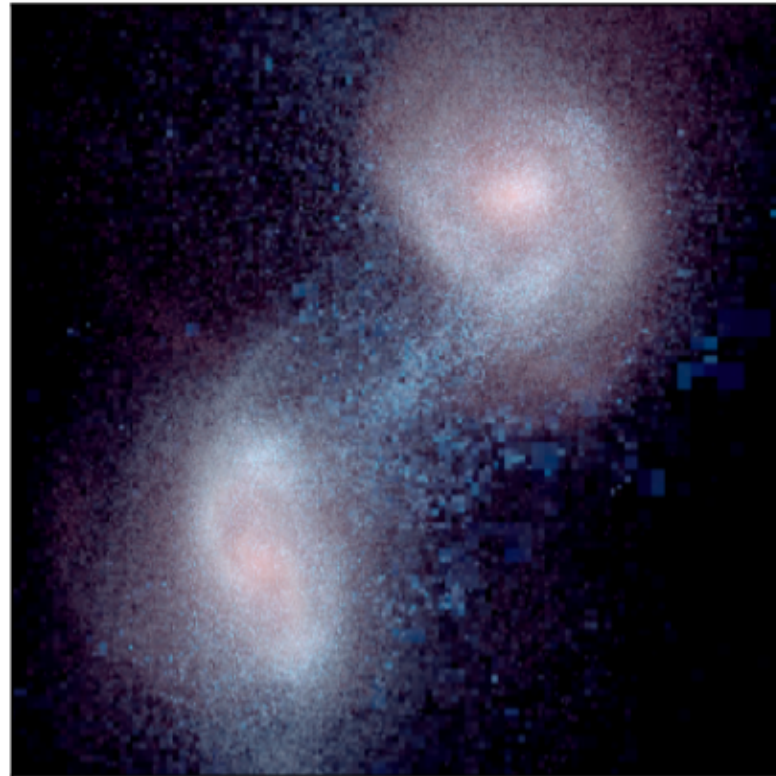
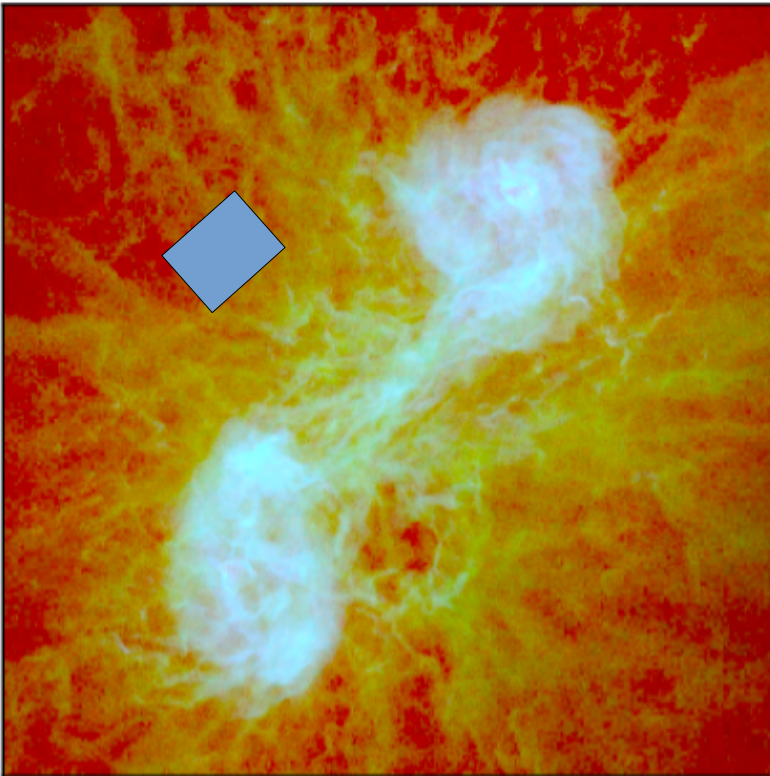


Future work



Next steps

- **Next steps: MHD, anisotropic thermal conduction.**
- **Link to cosmological simulations**



Conclusions

- **3D is necessary for a reliable treatment of bow shocks and the acceleration of cold clouds.**
- **Radiative cooling causes shattering of large clouds. Both in 2D and 3D. The gas in galaxy outskirts is probably in a different state than revealed by simulations.**
- **Next steps:
MHD, anisotropic conduction. Cosmology.**

