Max Planck Institute for Astrophysics



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### The (non-ionizing) interstellar radiation in resolved dwarf galaxies

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# KS relation in dwarf galaxies

Bigiel et al. 2010





Very steep slope ( $N=2\sim3$ ) below  $\Sigma_{\rm SFR} < 10 \, M_{\odot} \, \rm pc^{-2}$ Transition from  $t_{\rm dep} \sim 1 \, \rm Gyr$  to  $t_{\rm dep} \sim 100 \, \rm Gyr$  (extremely inefficient SF)

# The interstellar radiation field (ISRF)



#### What does ISRF do?





## Heating and cooling in MW

In the Milky Way (or typical spiral galaxies), the PE effect is the dominating heating mechanism.

While the metal line emission (CII, OI) dominates cooling.

The thermal balance in the ISM is largely controlled by these processes.





# Heating and cooling in MW

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The thermal balance in the ISM is largely controlled by these processes.

SN feedback drives turbulence and shocks which perturb the gas and cause scatter. However, the majority of gas is still in thermal equilibrium (i.e. they sit on the equil. curve)



### Modeling spatially varying ISRF in dwarf galaxies

Due to the low surface density and low dust abundance in dwarf galaxies, we can neglect dust extinction and calculate the ISRF directly from stars using the simple  $r^{-2}$  law.



To speed up, we calculate the ISRF using the tree approximation during the tree walk. Is it really a fair approximation? => need radiative transfer to confirm (T. Peter's talk).

# Simulations

#### **Physics:**

- gravity & hydrodynamics (Gadget-3 + modern SPH)
- non-equilibrium cooling & chemistry network (similar to SILCC)
- star formation (SFR =  $\epsilon \rho_{gas} / t_{ff}$ , threshold = 100 cm<sup>-3</sup> & 100 K,  $\epsilon$ =2%)
- SNII feedback (individual explosions)
- metal enrichment (SNII+AGB)

#### Setup:

- isolated dwarf galaxy
- $Z = 0.1Z_{\odot}$ , dust-to-gas ratio = 0.1%, uniform distribution
- SPH particle mass =  $4 M_{\odot}$  in order to resolve cold gas (no pressure floor)





 $M_{vir} = 2 \times 10^{10} M_{\odot}$ 

### Simulations



## Simulations



## Resolving individual SN feedback

In a medium density of 1 cm<sup>-3</sup> & Z=0.1Z $\odot$ , we need SPH particle mass ~ M $\odot$ .



### Numerical difficulties with too high resolution

Stellar population approach: every  $100 \text{ M}_{\odot}$  stellar population there is one SNII event.

What if your particle mass is below 100  $M_{\odot}$ ? (e.g.  $4M_{\odot}$ )

- stochastic injection: do one SNII feedback for every 25 particles.
- how about metal enrichment? There's only  $4M_{\odot}$  available...

Direct IMF sampling:

- assign an array of stellar masses  $m_i$  from IMF to a star particle until  $\sum_i m_i > m_{star}$
- the residual mass  $m_{res} = \sum_i m_i$   $m_{star}$  is borrowed from the next star particle

(  $\sum_{i} m_{i} > m_{res} + m_{star,2}$  for the next particle )

- if a star particle is assigned with zero mass, remove it



# Stellar masses



Massive stars die immediately

# The spatial distribution of ISRF



The majority of gas feels a smoothly varying ISRF.

Nearby the young stars, ISRF can be enhanced by orders of magnitude (bright PDRs).



## The spatial distribution of ISRF



### Radial trend of ISRF changes thermal balance



### Radial trend of ISRF changes thermal balance



### Radial trend of ISRF changes thermal balance



### But it doesn't seem to be relevant...

SN feedback creates turbulence and shocks fast enough to drive the gas out of thermal equilibrium



Cooling time  $t_{cool} \approx 50 \text{ Myr}$ when  $n=1 \text{ cm}^{-3} \& Z=0.1 Z_{\odot}$ => no time to cool...

SN feedback not just causes the scatter in the phase diagram, as is the case in spirals.

The PE heating becomes sub-dominant.

The balance between metal line cooling and the SN heating controls how much gas can cool and form stars.

### Even when PE heating is turned off completely

Without PE & CR heating, the equil. temperature should be zero... The only heating source here is the SN feedback.



### SN heating dominates PE heating

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star-forming gas

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# SFR and H<sub>2</sub> fraction



variable 
$$G_0$$
,  $\xi_{CR} = 3 \times 10^{-18} \text{ s}^{-1}$   
variable  $G_0$ ,  $\xi_{CR} = 0$   
 $G_0 = 1$ ,  $\xi_{CR} = 3 \times 10^{-18} \text{ s}^{-1}$   
 $G_0 = 0$ ,  $\xi_{CR} = 0$ 

ISRF only has a modest effect on SFR, even when it's switched off completely.

ISRF has a stronger effect on the  $H_2$  fraction. - strong ISRF destroys  $H_2$  in dense clouds.

Dense gas has very low H<sub>2</sub> fraction- H<sub>2</sub> is a bad tracer of star formation in dwarfs.

## KS relation

Agree well with observations with whatever PE and CR heating, as long as the SN feedback is on.



# In tension with Forbes+ 2016 ...?

Forbes+ 2016 concluded that PE heating alone can suppress SFR in dwarfs.



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# Why so different?

#### Differences:

ours	Forbes+2016
4M <sub>O</sub> x 100	2.5-10 pc
100 cm <sup>-3</sup> & 100 K	Jeans (~ $1-10$ cm <sup>-3</sup> )
2%	1%
3 Myr	individual
$1-10 \text{ M}_{\odot}/\text{pc}^2$	?
	ours 4M <sub>O</sub> x 100 100 cm <sup>-3</sup> & 100 K 2% 3 Myr 1-10 M <sub>O</sub> /pc <sup>2</sup>

# Summary

We conducted isolated dwarf galaxy simulations with detailed ISM physics and high enough resolution ( $m_{gas} = 4 M_{\odot}$ ) to resolve the cold and dense gas.

We implemented a spatially varying IRSF calculated directly from stars without free parameters, which naturally produces a smooth varying background ISRF, while nearby the sources the ISRF can get much higher.

The ISM in dwarfs is not in thermal equilibrium! SN feedback generates turbulence and shocks which dominates over PE heating, drives the gas out of thermal equilibrium and regulates star formation.

The ISRF has a significant effect on the  $H_2$  fraction, even though  $H_2$  is not relevant for star formation.

### Including ionizing radiation

Use the correct delay times of individual massive stars. Still much warmer than the equil. curves.



### No ionizing radiation



### Radial profile

